

Lecture 16

16.1

Let $Z = \frac{\partial}{\partial \bar{z}} + \frac{1}{2} \bar{z}$, $\bar{Z} = \frac{\partial}{\partial z} - \frac{1}{2} z$, where

$$\frac{\partial}{\partial z} = \frac{\partial}{\partial x} - i \frac{\partial}{\partial y}, \quad \frac{\partial}{\partial \bar{z}} = \frac{\partial}{\partial x} + i \frac{\partial}{\partial y}.$$

Theorem: For all z in \mathbb{C} ,

$$(Z e_{j,k})(z) = i(2k)^{1/2} e_{j,k-1}(z), \quad j=0,1,2,\dots, k=1,2,\dots,$$

and

$$(\bar{Z} e_{j,k})(z) = i(2k+2)^{1/2} e_{j,k+1}(z), \quad j,k=0,1,2,\dots.$$

Remarks: Z annihilates and \bar{Z} creates states in quantum mechanics.

Recall: We call L the twisted Laplacian with $\tau=1$.

Theorem: For all $j,k=0,1,2,\dots$,

$$L e_{j,k} = (2k+1) e_{j,k}$$

Theorem: Let $\tau \in \mathbb{R} \setminus \{0\}$. Then $\{e_{j,k}^\tau : j,k=0,1,2,\dots\}$ is an orthonormal basis for $L^2(\mathbb{C})$.

Proof: For nonnegative integers α, β, μ and ν

$$(e_{\alpha,\beta}^\tau, e_{\mu,\nu}^\tau) = |\tau| \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e_{\alpha,\beta}^\tau \left(\frac{\tau}{\sqrt{|\tau|}} q, \sqrt{|\tau|} p \right) \overline{e_{\mu,\nu}^\tau \left(\frac{\tau}{\sqrt{|\tau|}} q, \sqrt{|\tau|} p \right)} dq dp$$

for $j,k=0,1,2,\dots$.

Let $q' = \frac{\tau}{\sqrt{|\tau|}} q$ and $p' = \sqrt{|\tau|} p$. Then

$$\begin{aligned} & (e_{\alpha, \beta}^{\tau}, e_{\mu, \nu}^{\tau}) \\ &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e_{\alpha, \beta}^{\tau}(q', p') \overline{e_{\mu, \nu}^{\tau}(q', p')} dq' dp'. \end{aligned}$$

Therefore $\{e_{j, k}^{\tau} : j, k = 0, 1, 2, \dots\}$ is an orthonormal set in $L^2(\mathbb{C})$. Now, let $f \in L^2(\mathbb{C})$ be such that

$$(f, e_{j, k}^{\tau}) = 0, \quad j, k = 0, 1, 2, \dots$$

Then

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(q, p) \overline{e_{j, k}^{\tau}\left(\frac{\tau}{\sqrt{|\tau|}} q, \sqrt{|\tau|} p\right)} dq dp = 0$$

Let $q' = \frac{\tau}{\sqrt{|\tau|}} q$ and $p' = \sqrt{|\tau|} p$ be as above.

Then

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f\left(\frac{\sqrt{|\tau|}}{\tau} q', \frac{1}{\sqrt{|\tau|}} p'\right) \overline{e_{j, k}^{\tau}(q', p')} dq' dp' = 0$$

But $\{e_{j, k}^{\tau} : j, k = 0, 1, 2, \dots\}$ is an orthonormal basis for $L^2(\mathbb{C})$. Therefore

$$f\left(\frac{\sqrt{|\tau|}}{\tau} q', \frac{1}{\sqrt{|\tau|}} p'\right) = 0$$

for almost all $(q', p') \in \mathbb{R}^2$. Therefore $f(q', p') = 0$ for almost all (q', p') in \mathbb{C} . This completes the proof.

Complete Spectral Analysis of $L_\tau, \tau \in \mathbb{R} \setminus \{0\}$

Theorem: Let $\tau \in \mathbb{R} \setminus \{0\}$. Then for all $j, k = 0, 1, 2, \dots$,

$$L_\tau e_{j,k}^\tau = (2k+1)|\tau| e_{j,k}^\tau.$$

Remarks: For $j, k = 0, 1, 2, \dots$,

$e_{j,k}^\tau$ is an eigenfunction of L_τ corresponding to the eigenvalue $(2k+1)|\tau|$. So, there are infinitely many

eigenfunctions corresponding to the same eigenvalue $(2k+1)|\tau|$.

We say that $(2k+1)|\tau|$ is an eigenvalue of L_τ with infinite multiplicity.

The proof of the theorem is by a change of coordinates.

To be done next time (i.e., on Thursday, July 28, 2021).